

GLOBAL REGULARITY ON 3-DIMENSIONAL SOLVMANIFOLDS

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ABSTRACT. Let M be any 3-dimensional (nonabelian) compact solvmanifold. We apply the methods of representation theory to study the convergence of Fourier series of smooth global solutions to first order invariant partial differential equations $Df = g$ in $C^\infty(M)$. We show that smooth infinite-dimensional irreducible solutions, when they exist, satisfy estimates strong enough to guarantee uniform convergence of the irreducible (or primary) Fourier series to a smooth global solution.

1. INTRODUCTION

Let S be a solvable Lie group, and Γ a discrete subgroup of S with compact quotient $\Gamma \backslash S$. There is then a unique probability measure ν on $\Gamma \backslash S$ that is invariant under translation on the right by elements of S . The regular representation of S on $L^2(\Gamma \backslash S, \nu)$ decomposes into a direct sum of a countable number of irreducible unitary representations π of S , each of finite multiplicity m_π [G]. Let D be a *first order* differential operator with complex coefficients, left-invariant on S and viewed on $\Gamma \backslash S$. Let $(\Gamma \backslash S)^\wedge$ denote the dual object of $\Gamma \backslash S$. If $g \in C^\infty(\Gamma \backslash S)$ and if g_π is an orthogonal component of g corresponding to some irreducible unitary representation π , then $g_\pi \in C^\infty(\Gamma \backslash S)$ too [A-B]. Modulo unitary equivalence, we may think of g_π as being a C^∞ -vector in any concrete realization, or model, of π . We are interested in algebraically well-defined conditions on D under which the global solvability of $Df = g$ in $C^\infty(\Gamma \backslash S)$ is equivalent to the solvability of $\pi(D)f_\pi = g_\pi$ in the C^∞ -vectors for each π in the spectrum of $\Gamma \backslash S$. In a sense, we are looking for algebraic conditions on D for the reduction of a global (geometrical) problem on $\Gamma \backslash S$ to a collection of purely group (representation) theoretic problems, none of which needs to be regarded as living on the manifold $\Gamma \backslash S$. Informally speaking, operators D admitting such a reduction are called *globally regular* (Definition (1.1)).

In order to describe the results, we will recall the classical situation on a torus T^2 of two dimensions (the situation being similar for T^n with $n > 2$). Let $D = \alpha \partial / \partial x + \beta \partial / \partial y$ and suppose for simplicity that α and β are real.

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Then D is globally regular if and only if β/α is not a (transcendental) Liouville number. The problem with Liouville numbers is that, in solving for the Fourier transform of the solution function, *small divisors* occur. Now, every solvmanifold $\Gamma \backslash S$ contains the structure of a torus $T = \Gamma[S, S] \backslash S$ of dimension ≥ 1 , although this torus does not reflect any of the nonabelian structure of S . The only representations in $(\Gamma \backslash S)^\wedge$ which are not infinite dimensional are the one-dimensional characters of $\Gamma[S, S] \backslash S$. Since the presence of this torus is inescapable, we denote, for each $g \in C^\infty(\Gamma \backslash S)$, the sum of the one-dimensional components of g by g_0 . Then global regularity is defined as follows. Let $L^2(\Gamma \backslash S) = \bigoplus_{\pi \in (\Gamma \backslash S)^\wedge, j=1, \dots, m_\pi} H_{\pi, j}$ be any (noncanonical) irreducible decomposition of $L^2(\Gamma \backslash S)$.

(1.1) **Definition.** A left-invariant differential operator D on $\Gamma \backslash S$ is called globally regular if the three conditions

- (1) $g \in C^\infty(\Gamma \backslash S)$,
- (2) For each $\pi \in (\Gamma \backslash S)^\wedge$ and $\forall j = 1, \dots, m_\pi \exists$ a solution in $C^\infty(\Gamma \backslash S)$ to $Df_{\pi, j} = g_{\pi, j}$ (where $g_{\pi, j} = (\pi, j)$ -component of g), and
- (3) $\exists f_0 \in C^\infty(T)$ such that $Df_0 = g_0$,

imply that there exists a solution in $C^\infty(\Gamma \backslash S)$ to $Df = g$.

Note that the solutions in (2) could be found in any convenient realization of π .

In previous papers we have dealt with *nilpotent* S . On the simplest nilmanifolds, the 3-dimensional Heisenberg manifolds, every first order differential operator D in the complexified Lie algebra is globally regular [R2]. On more complicated nilmanifolds the problem of small divisors arises in the representation spaces of the group as well as on the associated torus. Moreover, if $D = X + iY$ is regular, both ad_X and ad_Y must map each step of the lower central series of the Lie algebra of S (nilpotent) onto a sufficiently large subset of the next step. The details are explained in [C-R1, p. 349]. The purpose of this paper is to investigate the global regularity of first order differential operators on 3-dimensional compact solvmanifolds. We show that, as in the case of the simplest nilmanifolds, every first order differential operator on a 3-dimensional compact solvmanifold is globally regular.

2. 3-DIMENSIONAL SOLVMANIFOLDS

All 3-dimensional compact solvmanifolds can be described (up to homeomorphism) as quotients of two groups S_h and S_r by their various cocompact discrete subgroups. The groups S_h and S_r can both be described as $R^2 \rtimes R^1$, where $(x, t)(x', t') = (x + A^t x', t + t')$. Here A^t is a 1-parameter subgroup of $SL(2, R)$ through a matrix $A \in SL(2, Z)$. S_h arises when the eigenvalues of A are $\lambda > 1$ and λ^{-1} , so that the orbits of R^1 in R^2 are *hyperbolic*. S_r arises when A^t is a compact group of *rotations* of R^2 . Let $N := R^2 \times \{0\}$, the (abelian) nilradical of S . The cocompact discrete subgroups Γ are described in [A-G-H], based upon the facts (due to Mostow) that $\Gamma \cap N$ is a discrete lattice in R^2 , and that the image of Γ under the natural projection $S \rightarrow S/N$ is a discrete lattice in R^1 . We remark that $\Gamma \backslash S_h$ is determined up to homeomorphism by the eigenvalue $\lambda > 1$ of A , and λ must be such that $\lambda + \lambda^{-1} \in Z$. For this reason, we denote the 'hyperbolic' manifolds $\Gamma_\lambda \backslash S_h$. Note however that

S_h is independent of the value of $\lambda > 1$. We have good use for the following lemma ((3.4) in [C-R2]).

(2.1) **Lemma.** *If $S_h = R^2 \rtimes R^1$ with the diagonalized matrix $A = \begin{pmatrix} \lambda & 0 \\ 0 & \lambda^{-1} \end{pmatrix}$, $\lambda > 1$, and if Γ_λ is a cocompact discrete subgroup of S_h , then $\Gamma_\lambda \cap N$ is an abelian lattice \mathcal{L} of points (α, β) having the property that the product $\alpha\beta$ is bounded away from zero, except of course at the identity.*

We need also the following

(2.2) **Corollary.** *In the setting of the lemma above, the dual lattice $\mathcal{L}^* = \{\chi_{(a,b)}: \mathcal{L} \rightarrow 1\}$ is also a lattice of points (a, b) such that the product ab is bounded away from 0 (except of course for $(a, b) = (0, 0)$).*

This corollary will give useful information about $(\Gamma \backslash S)_\infty$, the infinite dimensional representations in the spectrum of $\Gamma \backslash S$, in the hyperbolic case. For all 3-dimensional compact solvmanifolds, $(\Gamma \backslash S)_\infty$ is constructed as follows. Let $\chi_{(\alpha, \beta)} \in \hat{N}$, where $\chi_{(\alpha, \beta)}: \Gamma \cap N \rightarrow 1$. Now let M be the extension of N by the stabilizer of $\chi_{(\alpha, \beta)}$ in S , and extend χ to M , so $\pi_{\alpha\beta} := \text{Ind}_M^S(\chi_{(\alpha, \beta)}) \in \hat{S}$. If $S = S_h$ then $M = N$ and $M \backslash S \cong R$, whereas if $S = S_r$ then $M \cong N \times Z$ and $M \backslash S \cong R/Z = \text{the circle group}$. If $H_{(\chi, M)}$ is the standard Mackey induced representation space, then $H_{(\chi, M)} = \{f: S \rightarrow C | f(ms) = \chi_{(\alpha, \beta)}(m)f(s), |f| \in L^2(M \backslash S)\}$. Define $L: H_{(\chi, M)} \rightarrow L^2(\Gamma \backslash S)$ by $(Lf)(\Gamma s) = \sum_{\Gamma \cap M \backslash \Gamma} f(\gamma s)$. Then L is a right S -invariant injection. If $\text{Int}(\chi, M) = \{t \in R | \chi^{\exp tT}: \Gamma \cap M \rightarrow 1\}$ where $\chi^a(b) = \chi(a^{-1}ba)$ then the multiplicity of $\pi_{\alpha\beta}$ in $L^2(\Gamma \backslash S)$ equals the number of distinct Γ -orbits in $\text{Int}(\chi, M)$.

In the case of $S = S_h$, it is easiest to describe $\Gamma_\lambda \backslash S_h$ if we take $\Gamma_\lambda \cap N = Z^2$ and $A \in SL(2, Z)$. However, the model for $\pi_{\alpha\beta}$ is simplest if A is diagonalized ($A = \begin{pmatrix} \lambda & 0 \\ 0 & \lambda^{-1} \end{pmatrix}$ as in Lemma (2.1)) and then $\Gamma \cap N$ is more difficult to describe. Nevertheless, the Corollary (2.2) above shows that in this model $\{\pi_{\alpha\beta} \in (\Gamma_\lambda \backslash S_h)_\infty\}$ has $\alpha\beta$ bounded away from zero. We remark upon the fact that $\alpha\beta \neq 0$ prevents $\pi_{\alpha\beta}$ from being a representation of the (less well-behaved) ' $ax + b$ ' group which is a quotient of S_h .

We need to say a few words about the rotational three dimensional solvmanifolds $\Gamma \backslash S_r$ as well. Unlike the hyperbolic case, there are only finitely many $\Gamma \backslash S_r$ up to homeomorphism. We take $A \in SL(2, Z)$, but this time with *no* eigenvalues > 1 . Now A turns out to be similar to a *rotation* by $2\pi/p$, where $p = 1, 2, 3, 4$, or 6 (see [A-G-H]). Since S_r is independent of p (up to group isomorphism), we denote the distinct rotational three dimensional solvmanifolds by $\Gamma_p \backslash S_r$, $p = 1, 2, 3, 4$, or 6 .

Our main result is

(2.3) **Theorem.** *Let $\Gamma \backslash S$ be any nonabelian 3-dimensional compact solvmanifold. If $D \in \mathcal{S}^C$, the complexified Lie algebra of S , then D is globally regular on $\Gamma \backslash S$.*

Since this result has been proved in an earlier paper when S is nilpotent, we concern ourselves here only with the manifolds $\Gamma_\lambda \backslash S_h$ and $\Gamma_p \backslash S_r$. Note also that the associated torus T is 1-dimensional, so that small divisors *cannot*

occur on T . Thus condition (3) of Definition (1.1) will be satisfied automatically whenever (2) holds.

We will divide our proof of the theorem into two sections, one dealing with the rotational group S_r , the other with the hyperbolic one, S_h .

3. PROOF OF THE THEOREM—CASE OF $\Gamma_p \backslash S_r$

Let \mathcal{S}_r be the Lie algebra of S_r with a linear basis T, X, Y and the commutation relations $[T, X] = -\frac{2\pi}{p}Y$ and $[T, Y] = \frac{2\pi}{p}X$. Let $\pi_{\alpha\beta}$ be a generic infinite-dimensional representation in $(\Gamma_p \backslash S_r)^*$, acting in $L^2(T)$, where $T = M \backslash S_r$, a 1-dimensional torus as described in §2. For $f \in L^2(T)$ the action of $\pi_{\alpha\beta}$ on f is given by

$$\pi_{\alpha\beta}(x, y; t) f(\tau) = \exp(2\pi i \langle (\alpha, \beta) \sigma(\frac{2\pi}{p}\tau), (x, y) \rangle) f(\tau + t),$$

where $\sigma(s) = \begin{pmatrix} \cos s & \sin s \\ -\sin s & \cos s \end{pmatrix}$. For the basis vector fields X, Y, T this amounts to

$$\begin{aligned} d\pi_{\alpha\beta}(X) &= 2\pi i \left(\alpha \cos \frac{2\pi}{p}t - \beta \sin \frac{2\pi}{p}t \right), \\ d\pi_{\alpha\beta}(Y) &= 2\pi i \left(\alpha \sin \frac{2\pi}{p}t + \beta \cos \frac{2\pi}{p}t \right), \\ d\pi_{\alpha\beta}(T) &= \frac{d}{dt}. \end{aligned} \quad (3.1)$$

Since the constant p plays a negligible role in the proof (even though it classifies the rotational 3-dimensional solvmanifolds), we will set $p = 1$ in what follows.

We will break the proof into two cases of $D \in \mathcal{S}_r^C$.

Case 1. $D = X + \gamma Y$, $\gamma \in C$. This essentially covers all $D \in \mathcal{N}^C$ (up to a constant factor) since $D = Y$ and $D = X$ behave alike and the case of $D = 0$ is trivial.

Case 2. $D = T + i(aX + bY)$, $a, b \in R$. This covers all $D \in \mathcal{S}_r^C \sim \mathcal{N}^C$ (up to an isomorphism) because the real part of D can be absorbed into T .

Proof of Case 1. Write $D = X + (a + ib)Y$, $a, b \in R$. Then, in view of (3.1), the operator $d\pi_{\alpha\beta}(D)$ is a multiplication by the function

$$\begin{aligned} D_{\alpha\beta}(t) &:= 2\pi i[(\alpha + a\beta + ib\beta) \cos 2\pi t + (a\alpha - \beta + ib\alpha) \sin 2\pi t] \\ &= 2\pi i(w \cos 2\pi t + z \sin 2\pi t). \end{aligned} \quad (3.2)$$

By hypothesis, the equation $d\pi_{\alpha\beta}(D)f_{\alpha\beta} = g_{\alpha\beta}$ has the solution

$$f_{\alpha\beta}(t) = g_{\alpha\beta}(t)/D_{\alpha\beta}(t). \quad (3.3)$$

To prove the theorem we need to show that $\sum_{(\alpha, \beta)} Lf_{\alpha\beta} \in C^\infty(\Gamma \backslash S)$. Here (α, β) varies over a cross section of Γ -orbits, so that each infinite-dimensional primary summand will be spanned. By the Auslander-Brezin version of the Sobolev inequality it suffices to show $\sum_{(\alpha, \beta)} \|ULf_{\alpha\beta}\|^2 < \infty$ for all $U \in \mathcal{U}(\mathcal{S}_r)$. Since L is an S -invariant isometry from $H_{\alpha\beta}$ into $L^2(\Gamma \backslash S)$, this is the same as to show $\sum_{(\alpha, \beta)} \|Uf_{\alpha\beta}\|_2^2 < \infty$. We begin by estimating the sum

$\sum \|f_{\alpha\beta}\|_2^2$. The problem, of course, is that $D_{\alpha\beta}(t)$ in (3.3) can have zeros, and that even when it has no zeros, we must know how close $|D_{\alpha\beta}(t)|$ can come to 0 as $(\alpha, \beta) \rightarrow \infty$. We can write

$$D_{\alpha\beta}(t) = \pi i((w - iz)e^{2\pi it} + (w + iz)e^{-2\pi it}) = \pi i(Ae^{2\pi it} + Be^{-2\pi it}),$$

with A and B defined by the last equation. The minimum of $|D_{\alpha\beta}(t)|$ occurs where A and B are rotated to opposite directions, and then

$$\begin{aligned} \text{Min } |D_{\alpha\beta}(t)| &= \pm\pi(|A| - |B|) \\ (3.4) \quad &= \pi(\alpha^2 + \beta^2)^{1/2}((a^2 + (1+b)^2)^{1/2} - (a^2 + (1-b)^2)^{1/2}) \\ &= (\alpha^2 + \beta^2)^{1/2} \cdot K \geq \sqrt{2}|\alpha\beta|^{1/2} \cdot K. \end{aligned}$$

If $b \neq 0$, the constant K is $\neq 0$. Since $|\alpha\beta| \gg 0$, $|f_{\alpha\beta}(t)| \leq C|g_{\alpha\beta}(t)|$ with the constant C independent of (α, β) . Consequently, $\sum_{(\alpha, \beta)} \|f_{\alpha\beta}\|^2 \leq C \sum_{(\alpha, \beta)} \|g_{\alpha\beta}\|^2 < \infty$ for $g \in C^\infty(\Gamma \backslash S)$.

If $b = 0$, $D_{\alpha\beta}(t)$ does have one or more zeros and

$$D_{\alpha\beta}(t) = 2\pi i((\alpha + a\beta) \cos 2\pi t + (a\alpha - \beta) \sin 2\pi t).$$

If $D_{\alpha\beta}(t_0) = 0$, then $g_{\alpha\beta}(t_0) = 0$ too, since $f_{\alpha\beta}$ given by (3.3) is in $C^\infty(T)$. The idea is to control $f_{\alpha\beta}$ inside specified intervals around each of the t_0 's by the mean value theorem, and to use the monotonicity of $D_{\alpha\beta}$ on large exterior intervals to control $|f_{\alpha\beta}|$ by keeping $|D_{\alpha\beta}|$ big. We have

$$D'_{\alpha\beta}(t_0) = -2\pi(\alpha^2 + \beta^2)^{1/2}(1 + a^2)^{1/2},$$

and by the mean value theorem

$$|D'_{\alpha\beta}(t) - D'_{\alpha\beta}(t_0)| \leq |t - t_0|4\pi^2(\alpha^2 + \beta^2)^{1/2}(1 + a^2)^{1/2}.$$

Consequently, for $|t - t_0| < 1/4\pi$, $|D'_{\alpha\beta}(t)| \geq \pi(\alpha^2 + \beta^2)^{1/2}(1 + a^2)^{1/2}$ and, since $D_{\alpha\beta}(t_0) = 0$,

$$|D_{\alpha\beta}(t)| \geq \pi(\alpha^2 + \beta^2)^{1/2}(1 + a^2)^{1/2}|t - t_0|$$

again by the mean value theorem. For $g_{\alpha\beta}(t)$ we have the estimate

$$(3.5) \quad |g_{\alpha\beta}(t)| = |g_{\alpha\beta}(t) - g_{\alpha\beta}(t_0)| \leq \|Tg_{\alpha\beta}\|_\infty |t - t_0|,$$

$\|\cdot\|_\infty$ denoting the sup norm on the torus $M \backslash S$. So

$$(3.6) \quad |f_{\alpha\beta}(t)| \leq \|Tg_{\alpha\beta}\|_\infty \frac{1}{\pi}(\alpha^2 + \beta^2)^{-1/2}(1 + a^2)^{-1/2}, \quad \text{for } |t - t_0| < 1/4\pi.$$

On the intervals complementary to $|t - t_0| < 1/4\pi$,

$$|D_{\alpha\beta}(t)| \geq |D_{\alpha\beta}(t_0 \pm 1/4\pi)| \geq \pi(\alpha^2 + \beta^2)^{1/2}(1 + a^2)^{1/2}/4\pi,$$

hence

$$(3.7) \quad |f_{\alpha\beta}(t)| \leq 4\|g_{\alpha\beta}\|_\infty(\alpha^2 + \beta^2)^{-1/2}(1 + a^2)^{-1/2} \quad \text{for } |t - t_0| \geq 1/4\pi$$

for each of the two values t_0 of t where $D_{\alpha\beta}$ vanishes. By (3.6) and (3.7), for all t on the torus $M \backslash S$ we have the following estimate:

$$(3.8) \quad |f_{\alpha\beta}(t)| \leq (\|Tg_{\alpha\beta}\|_\infty/\pi + 4\|g_{\alpha\beta}\|_\infty)(\alpha^2 + \beta^2)^{-1/2}(1 + a^2)^{-1/2}.$$

By Sobolev's inequality we may replace the sup norms in (3.8) by L^2 -norms of $g_{\alpha\beta}$, $Tg_{\alpha\beta}$, and $T^2g_{\alpha\beta}$. The sum $\sum_{(\alpha, \beta)} \|f_{\alpha\beta}\|_2^2$ is finite because each sum $\sum_{(\alpha, \beta)} \|T^k g_{\alpha\beta}\|_2^2$ for $k = 0, 1, 2$ is finite and $|\alpha\beta|$ is bounded away from zero.

Next, we must show $\sum_{(\alpha, \beta)} \|U f_{\alpha\beta}\|_2^2 < \infty$ for every fixed $U \in \mathcal{U}(\mathcal{S}_r)$. Since $[D, \mathcal{N}] = 0$, this is true for all $U \in \mathcal{U}(\mathcal{N})$. It remains to show $\sum_{(\alpha, \beta)} \|T^k f_{\alpha\beta}\|_2^2 < \infty$ for $k = 1, 2, \dots$, because every $U \in \mathcal{U}(\mathcal{S}_r)$ can be written as a linear combination of monomials $T^k V$ with V in $\mathcal{U}(\mathcal{N})$.

(3.9) Proposition. For $D = X + \gamma Y$, $\gamma \in C$ and $k = 1, 2, 3, \dots$ the $(k+1)$ -fold bracket product $[D[D \cdots [D, T^k] \cdots]]$ is 0.

Proof of Proposition. By Leibnitz's rule for the derivation $[D, \cdot]$ of the algebra $\mathcal{U}(\mathcal{S}_r)$ we have $[D, T^k] = \sum_{j=1}^k T \cdots T [D, T] T \cdots T$, with $[D, T]$ at j th place. Since D commutes with \mathcal{N} and $[D, T] = -X + \gamma Y \in \mathcal{N}$, the derivation repeated $k+1$ times is zero.

(3.10) Proposition. For $f_{\alpha\beta}$ as in (3.3) we have

$$(3.11) \quad T^k f_{\alpha\beta} = h_k / D_{\alpha\beta}^{k+1}$$

with

$$h_k = [D[D \cdots [D, T^k] \cdots]] g_{\alpha\beta} + D[D \cdots [D, T^k] \cdots] g_{\alpha\beta} \\ + \cdots + D^{k-1} [D, T^k] + D^k T^k g_{\alpha\beta}.$$

The first bracket involves k D 's with the number of D 's inside the brackets decreasing by one in each successive summand.

Proof of Proposition. In view of Proposition (3.9) this is formula (1.8) on p. 353 of [C-R1].

The estimates on $T^k f_{\alpha\beta}$ given by (3.11) can now be done in a manner similar to that already presented. If $D_{\alpha\beta} \neq 0$, we use the inequality (3.4) raised to the power $k+1$ to estimate the denominator $|D_{\alpha\beta}|^{k+1}$. If $D_{\alpha\beta}(t_0) = 0$, the numerator h_k in (3.11) must have a $(k+1)$ th order zero at t_0 since $T^k f_{\alpha\beta}$ is C^∞ . Instead of the estimate (3.5) we use

$$|h_k(t)| \leq \|T^{k+1} h_k\|_\infty |t - t_0|^{k+1} / (k+1)!$$

which follows from Taylor's formula. In the denominator of (3.6) and (3.7) we use the $(k+1)$ th power of the previous estimate for $D_{\alpha\beta}$.

Proof of Case 2. Here $D = T + i(aX + bY)$, $a, b \in R$, and

$$(3.12) \quad d\pi_{\alpha\beta}(D) = \frac{d}{dt} - 2\pi((a\alpha + b\beta) \cos 2\pi t + (b\alpha - a\beta) \sin 2\pi t) \\ = \frac{d}{dt} - 2\pi(A \cos 2\pi t + B \sin 2\pi t).$$

Write $A \cos 2\pi t + B \sin 2\pi t = (A^2 + B^2)^{1/2} \sin 2\pi(t + \phi)$, with the constant ϕ depending upon A and B . If $d\pi_{\alpha\beta}(D) f_{\alpha\beta} = g_{\alpha\beta}$, we have

$$f_{\alpha\beta}(t) = \exp(-2\pi(A^2 + B^2)^{1/2} \cos 2\pi(t + \phi)) \\ \times \left(\int_{-\phi-1/2}^t g_{\alpha\beta}(\tau) \exp(2\pi(A^2 + B^2)^{1/2} \cos 2\pi(\tau + \phi)) d\tau + C \right).$$

Here we identify $M \setminus S = R/Z$ with the interval $[-\phi - \frac{1}{2}, \frac{1}{2} - \phi]$, and we use the fact that $f_{\alpha\beta}$, $g_{\alpha\beta}$, and the exponentials all have period 1. Since C is

arbitrary, choosing $C = 0$ and changing the variables $\tau' = \tau + \phi$, $t' = t + \phi$ we have

$$\begin{aligned} f_{\alpha\beta}(t - \phi) &= \int_{-1/2}^t g(\tau - \phi) \exp(2\pi(A^2 + B^2)^{1/2}(\cos 2\pi\tau - \cos 2\pi t)) d\tau \\ &= - \int_t^{1/2} (\text{same integrand as above}). \end{aligned}$$

The last equality follows from the periodicity of $f_{\alpha\beta}$, because $0 = f_{\alpha\beta}(-1/2 - \phi) = f_{\alpha\beta}(1/2 - \phi)$. For the estimates on $f_{\alpha\beta}(t)$ or on $f_{\alpha\beta}(t - \phi)$ we use the integral over $-1/2 \leq \tau \leq t$ if $-1/2 \leq t \leq 0$ and over $t \leq \tau \leq 1/2$ for $0 < t < 1/2$. That way we always have $\cos 2\pi\tau < \cos 2\pi t$, so that the exponent in the integral defining $f_{\alpha\beta}$ is negative. Then

$$\|f_{\alpha\beta}\|_{\infty} \leq C \|g_{\alpha\beta}\|_{\infty} \leq C_1 (\|g_{\alpha\beta}\|_2 + \|T g_{\alpha\beta}\|_2)$$

Since the $f_{\alpha\beta}$'s all live on an interval of length 1, the same inequality holds for $\|f_{\alpha\beta}\|_2$ with a constant, say C'_1 , independent of (α, β) . Thus $\sum_{(\alpha, \beta)} \|f_{\alpha\beta}\|_2^2 < \infty$ and $f = \sum f_{\alpha\beta} \in L^2(\Gamma_p \backslash S_r)$. Next, to estimate $\sum_{(\alpha, \beta)} \|U f_{\alpha\beta}\|_2^2$ for every $U \in \mathcal{U}(\mathcal{S}_r)$ it suffices to consider $U = X^k Y^j T^l$. By (3.12) T just differentiates $f_{\alpha\beta}$ yielding $2\pi(A^2 + B^2)^{1/2} \sin 2\pi(t + \phi) f_{\alpha\beta}(t) + g_{\alpha\beta}(t)$ with A, B depending linearly on α, β . Successive powers of T differentiate $g_{\alpha\beta}$ and $\sin 2\pi(t + \phi)$, or $f_{\alpha\beta}(t)$. Operating by X 's or Y 's just multiplies by a polynomial in (α, β) times a (bounded) combination of sines and cosines. However, $X^2 + Y^2$ acts on $L^2(M \backslash S)$ by multiplying by $-4\pi^2(\alpha^2 + \beta^2)$, so that $\sum_{(\alpha, \beta)} (\alpha^2 + \beta^2)^q \|g_{\alpha\beta}\|_2^2 < \infty$. Thus $\sum_{(\alpha, \beta)} \|X^k Y^j T^l f_{\alpha\beta}\|_2^2$ can be estimated by $C \sum \|(X^2 + Y^2)^q g_{\alpha\beta}\|_2^2$ for some q .

Remark. Although we have pretended $p = 1$ in the arguments above, this was just a notational convenience, and the same arguments work just as well for general $\Gamma \backslash S_r$. Actually, the use of $\Gamma_1 \backslash S_r$ has some independent interest however. Here, although S_r is not nilpotent, $\Gamma \cong \mathbb{Z}^3$. Thus, by the general theory of Mostow [M], $\Gamma_1 \backslash S_r$ is homeomorphic to a torus T^3 . Here we get a proof of global regularity for elements of \mathcal{S}_r^C acting on T^3 , keeping in mind that we mean global regularity with respect to the representation theory of S_r . Thus we have a natural class of operators on T^3 which are best analyzed with respect to S_r .

4. PROOF OF THE THEOREM—CASE OF $\Gamma_\lambda \backslash S_h$

Let \mathcal{S}_h be the Lie algebra of S_h with a linear basis T, X, Y and the commutation relations $[T, X] = X \ln \lambda$ and $[T, Y] = -Y \ln \lambda$. A generic infinite dimensional representation $\pi_{\alpha\beta}$ in $(\Gamma_\lambda \backslash S_h)^\wedge$ acts on $L^2(R)$ by

$$\pi_{\alpha\beta}(x, y, t)f(\tau) = \exp 2\pi i(\alpha\lambda^\tau x + \beta\lambda^{-\tau} y)f(\tau + t).$$

For the basis vector fields X, Y, T this amounts to

$$(4.1) \quad d\pi_{\alpha\beta}(X) = 2\pi i\alpha\lambda^t, \quad d\pi_{\alpha\beta}(Y) = 2\pi i\beta\lambda^{-t}, \quad d\pi_{\alpha\beta}(T) = \frac{d}{dt}.$$

The space $H_{\alpha\beta}^\infty$ of C^∞ -vectors for $\pi_{\alpha\beta}$ consists of $C^\infty(R)$ functions ϕ with

$$\lim_{|t| \rightarrow \infty} \lambda^{m|t|} \phi^{(n)}(t) = 0$$

for every $m, n \in N$. These ϕ are called “super-Schwartz”. We will denote this space by $\mathcal{S}(R)$.

We will break the proof into four cases of $D \in \mathcal{S}_h^C$.

Case 1. $D = T$. Any nonzero complex multiple of $D = T + aX + bY \in \mathcal{S}_h$ can be reduced up to isomorphism to $D = T$, if $a, b \in R$.

Case 2. $D = aX + bY \in \mathcal{N} \subset \mathcal{S}_h$, $a, b \in R$.

Case 3. $D = aX + bY \in \mathcal{N}^C \subset \mathcal{S}_h^C$, $a, b \in C$.

Case 4. $D = T + i(aX + bY)$, $a, b \in R$.

Proof of Case 1. $D = T$. The equation $d\pi_{\alpha\beta}(T)f_{\alpha\beta} = g_{\alpha\beta}$ becomes $(d/dt)f_{\alpha\beta} = g_{\alpha\beta}$. Since both $f_{\alpha\beta}$ and $g_{\alpha\beta}$ are in $H_{\alpha\beta}^\infty$, we have

$$(4.2) \quad f_{\alpha\beta}(t) = \int_{-\infty}^t g_{\alpha\beta}(\tau) d\tau = - \int_t^\infty g_{\alpha\beta}(\tau) d\tau.$$

Subcase 1a. If $t \geq 0$, we use the second integral for the estimates on $f_{\alpha\beta}$ and the fact that $g_{\alpha\beta}(\tau) = Xg_{\alpha\beta}(\tau)/2\pi i\alpha\lambda^\tau$, so that

$$(4.3) \quad |f_{\alpha\beta}(t)| \leq \|Xg_{\alpha\beta}\|_\infty (2\pi\alpha \ln \lambda)^{-1} \lambda^{-t} \in L^2([0, \infty), dt).$$

Squaring, integrating $\int_0^\infty \cdots dt$, and applying Sobolev’s inequality to estimate the norm $\|Xg_{\alpha\beta}\|_\infty$ by a combination of L^2 -norms of derivatives of $Xg_{\alpha\beta}$ we obtain

$$(4.4) \quad \|f_{\alpha\beta}\|_{L^2[0, \infty)}^2 \leq \frac{c}{\alpha} (\|Xg_{\alpha\beta}\|_2^2 + \|TXg_{\alpha\beta}\|_2^2),$$

for some constant c . As in the proof of Lemma (2.1) (see [C-R2]), we pick the representative (α, β) from the orbit corresponding to π so that $\lambda^{-2} \leq |\beta/\alpha| \leq 1$. Since $|\alpha\beta| \gg 0$, we must have $|\alpha| \gg 0$ too. Thus $\sum_{(\alpha, \beta)} \|f_{\alpha\beta}\|_{L^2[0, \infty)}^2 < \infty$.

Subcase 1b. If $t < 0$, we use the first integral formula for $f_{\alpha\beta}$ in (4.2), together with the fact that $g_{\alpha\beta}(\tau) = Yg_{\alpha\beta}(\tau)/2\pi i\beta\lambda^{-\tau}$. Picking (α, β) as in Subcase 1a, we insure that $|\beta| \gg 0$. Now an argument almost identical to that in Subcase 1a shows that $f = \sum f_{\alpha\beta} \in L^2(R)$.

To complete Case 1, we need to prove that $\sum \|Uf_{\alpha\beta}\|_2^2 < \infty$, for each fixed $U \in \mathcal{U}(\mathcal{S}_h)$. If $U = T^p$ there is no problem, since $[D, U] = 0$. If $U = X^m Y^n T^p$, $Uf_{\alpha\beta}(t) = (2\pi i)^{m+n} \alpha^m \beta^n \lambda^{(m-n)t} T^{p-1} g_{\alpha\beta}(t)$. If $p \geq 1$, there is no problem. But if $p = 0$, we estimate $|g_{\alpha\beta}(\tau)|$ by a $\|Vg_{\alpha\beta}\| \lambda^{kt}$, where $V \in \mathcal{U}(\mathcal{S}_h)$ and $k \in \mathbb{Z}$ are chosen one way for $t < 0$ and another way for $t \geq 0$, so as to assure that $Uf_{\alpha\beta} \in L^2(-\infty, 0)$ and $Uf_{\alpha\beta} \in L^2(0, \infty)$ separately. Hence $\sum Uf_{\alpha\beta} \in L^2(R)$.

Case 2. Suppose $D = aX + bY$, with a and b real. Here,

$$f_{\alpha\beta}(t) = g_{\alpha\beta}(t)/2\pi i(a\alpha\lambda^t + b\beta\lambda^{-t}).$$

We will need to consider the occurrence of zeros in the denominator. Observe that the denominator vanishes if and only if $\lambda^{2t} = -b\beta/a\alpha$, which means $t = t_0 = \frac{1}{2} \log_\lambda(-b\beta/a\alpha)$.

Subcase 2a. Suppose $b\beta/a\alpha > 0$, so that $h(t) = (a\alpha\lambda^t + b\beta\lambda^{-t})^{-1}$ has a maximum (absolute) value on R . At the maximum point $h'(t) = 0$, so that

$\lambda^{2t} = b\beta/a\alpha$. Thus $h(t_{\max}) = \lambda^{t_{\max}}/2b\beta = \pm 1/2|ab\alpha\beta|^{1/2}$. Thus $|f_{\alpha\beta}(t)| \leq |g_{\alpha\beta}(t)|/2|ab\alpha\beta|^{1/2}$. Since $|\alpha\beta| \gg 0$, we have $\sum f_{\alpha\beta} \in L^2$, where the sum is over those (α, β) such that $b\beta/a\alpha > 0$. The same is true for $\sum Uf_{\alpha\beta}$, for each fixed $U \in \mathcal{U}(\mathcal{S}_h)$ since

- (1) if $U \in \mathcal{U}(\mathcal{N})$, $[U, D] = 0$, and
- (2) if $U = X^m Y^n T^k$, then

$$\frac{d}{dt}(g_{\alpha\beta}(t)h(t)) = g'_{\alpha\beta}h + g_{\alpha\beta}h' = (Tg_{\alpha\beta})h + g_{\alpha\beta}h',$$

and $g_{\alpha\beta}h' = (X - Y)g_{\alpha\beta} \cdot h^2_{\alpha\beta}$, which is still L^2 -summable.

Subcase 2b. Either $a\alpha = 0$ or $b\beta = 0$. Since $|\alpha\beta| \gg 0$ (except at $(0, 0)$) this means $a = 0$ or $b = 0$ (but not both). If $b = 0$, then $f_{\alpha\beta}(t) = -Yg_{\alpha\beta}(t)/4\pi^2 a\alpha\beta$. Then $\sum f_{\alpha\beta} \in L^2$, since $|\alpha\beta| \gg 0$. A similar argument applies if $a = 0$.

Subcase 2c. Here, we suppose there exists a zero, $t_0 = \frac{1}{2}\log_\lambda(-b\beta/a\alpha)$. Then $g_{\alpha\beta}(t_0) = 0$ too, and $f_{\alpha\beta}(t) = (g_{\alpha\beta}(t) - g_{\alpha\beta}(t_0))/2\pi i(a\alpha\lambda^t + b\beta\lambda^{-t})$. Since it would suffice to prove separately the summability of the real and imaginary parts of $f_{\alpha\beta}$, we can assume wlog that $g_{\alpha\beta}$ is real-valued. For some τ between 0 and t , we have $f_{\alpha\beta}(t) = g'_{\alpha\beta}(\tau)/2\pi i(a\alpha\lambda^\tau - b\beta\lambda^{-\tau})\ln\lambda$, by the general mean value theorem. Since $b\beta/a\alpha < 0$ in this case, the new denominator cannot vanish. Hence, recalling that τ depends on t ,

$$\begin{aligned} 4\pi^2 \int_R |f_{\alpha\beta}(t)|^2 dt &= \int_{|t-t_0| \leq 1} |g'_{\alpha\beta}(\tau)|^2 |a\alpha\lambda^\tau - b\beta\lambda^{-\tau}|^{-2} |\ln\lambda|^{-2} dt \\ &\quad + \int_{|t-t_0| \geq 1} |g_{\alpha\beta}(t)|^2 |a\alpha\lambda^t + b\beta\lambda^{-t}|^{-2} dt \\ &\leq \frac{\|g'_{\alpha\beta}\|_\infty^2}{-4a\alpha b\beta \ln^2 \lambda} + \int_{|t-t_0| \geq 1} |g_{\alpha\beta}(t)|^2 \cdot \text{Max} |a\alpha\lambda^t + b\beta\lambda^{-t}|^{-2} dt. \end{aligned}$$

The first summand is summable over (α, β) , by Sobolev's inequality. However, by the mean value theorem, $|a\alpha\lambda^t + b\beta\lambda^{-t}| = |(a\alpha\lambda^\tau - b\beta\lambda^{-\tau})(\ln\lambda)(t - t_0)| \geq 2(-a\alpha b\beta)^{1/2} |\ln\lambda|$, if $|t - t_0| \geq 1$. Thus

$$\pi^2 \int_R |f_{\alpha\beta}(t)|^2 dt \leq \frac{\|g'_{\alpha\beta}\|_\infty^2}{|a\alpha b\beta| \ln^2 \lambda} + \int_R \frac{|g_{\alpha\beta}(t)|^2}{|a\alpha b\beta| \ln^2 \lambda} dt,$$

and so $\sum f_{\alpha\beta} \in L^2$.

If $U \in \mathcal{U}(\mathcal{N})$, then $\sum Uf_{\alpha\beta} \in L^2$ too, since $[D, U] = 0$. So it suffices to consider $\sum T^k f_{\alpha\beta}$, for each fixed $k \in \mathbb{Z}^+$. First, consider $k = 1$. $2\pi i T f_{\alpha\beta}(t) = h(t)(a\alpha\lambda^t + b\beta\lambda^{-t})^{-2}$, where $h(t) = (a\alpha\lambda^t + b\beta\lambda^{-t})g'_{\alpha\beta}(t) - g_{\alpha\beta}(t)(a\alpha\lambda^t - b\beta\lambda^{-t})\ln\lambda$. Here $h(t)$ must have a zero of order at least two at $t = t_0$. Therefore, using a Taylor remainder of degree two in $(t - t_0)$, there exist τ and τ' such that $T f_{\alpha\beta}(t) = h''(\tau')/D''(\tau)2\pi i$ where $D(t) = (a\alpha\lambda^t + b\beta\lambda^{-t})\ln\lambda$. But

$$|D''(t)| = 2[(a\alpha\lambda^\tau - b\beta\lambda^{-\tau})\ln\lambda]^2 + [(a\alpha\lambda^\tau + b\beta\lambda^{-\tau})\ln\lambda]^2 \geq 8|a\alpha b\beta| \ln^2 \lambda,$$

since the second square is positive. Thus $|Tf_{\alpha\beta}(t)| \leq |h''(\tau')|/16\pi|a\alpha b\beta|$. Hence

$$4\pi^2 \int_R |Tf_{\alpha\beta}|^2 \leq \int_{|t-t_0| \leq 1} \|h''\|_\infty / 64(a\alpha b\beta)^2 dt + \int_R |h(t)|^2 / \text{Min}\{|D(t)|^2 \mid |t-t_0| \geq 1\} dt.$$

But $\|h''\|_\infty$ is bounded by norms of derivatives of $g_{\alpha\beta}$, providing good summability over (α, β) of the first term. Also, $|D(t)|^2 = |D''(\tau)(t-t_0)^2/2!|^2 \geq 16(a\alpha b\beta)^2 \ln^4 \lambda \gg 0$. Thus $\sum Tf_{\alpha\beta} \in L^2$.

Next, we consider $k > 1$. We make the following observations.

$$(4.5) \quad [D, T^k] = \sum_{j=1}^k T \cdots T[D, T]T \cdots T,$$

where $[D, T]$ is in the j th position. Hence

$$(4.6) \quad \begin{aligned} [D, [D, T^k]] &= \sum_{i,j=1}^k T \cdots T[D, T]T \cdots T \\ &= 2 \sum_{1 \leq i < j \leq k} T \cdots T[D, T]T \cdots T[D, T]T \cdots T, \end{aligned}$$

where $[D, T]$ occupies the i th and j th positions. (This is a result of \mathcal{N} being abelian and normal.) The k -fold bracket

$$(4.7) \quad [D, [D, \dots [D, T^k] \cdots]] = k! [D, T]^k,$$

while the $(k+1)$ -fold bracket vanishes. By (1.8) of [C-R1, p.353],

$$(4.8) \quad 2\pi iT^k f_{\alpha\beta}(t) = h_k(a\alpha\lambda^t + b\beta\lambda^{-t})^{-(k+1)}$$

where

$$\begin{aligned} h_k &= [D[D \cdots [D, T^k] \cdots]]g_{\alpha\beta} + D[D \cdots [D, T^k] \cdots]g_{\alpha\beta} \\ &\quad + \cdots + D^{k-1}[D, T^k]g_{\alpha\beta} + D^k T^k g_{\alpha\beta}. \end{aligned}$$

It follows that $\|T^k f_{\alpha\beta}\|_2^2 = \int_{|t-t_0| \leq 1} + \int_{|t-t_0| > 1}$, with the integrand being determined by (4.8).

The second integral can be estimated as in the case of $k = 1$, while the denominator in the integrand in the first integral has a Taylor expansion in $(t-t_0)$ using derivatives of order $\leq k+1$. For $k > 1$ and odd, the $(k+1)$ th derivative of $(a\alpha\lambda^t + b\beta\lambda^{-t})^{k+1}$ is of the form $(d/dt)^{k+1}[(a\alpha\lambda^t + b\beta\lambda^{-t})^{k+1}] = \sum_{j=0}^{(k+1)/2} c_j(a\alpha\lambda^t + b\beta\lambda^{-t})^{k+1-2j}(a\alpha\lambda^t - b\beta\lambda^{-t})^{2j} \ln^k \lambda$ with all $c_j \geq 0$ and $c_{(k+1)/2} > 0$. Thus, if $k+1$ is even,

$$\frac{d^{k+1}}{dt^{k+1}} (a\alpha\lambda^t + b\beta\lambda^{-t})^{k+1} \geq C(a\alpha\lambda^t - b\beta\lambda^{-t})^{k+1}, \quad \text{where } C > 0.$$

If $k+1$ is odd, then

$$\frac{d^{k+1}}{dt^{k+1}} (a\alpha\lambda^t + b\beta\lambda^{-t})^{k+1} \geq C'|a\alpha\lambda^t - b\beta\lambda^{-t}|^{k+1}, \quad \text{for } C' > 0,$$

by a similar calculation. Hence the minima over $|t - t_0| \leq 1$ in the resulting estimates proceed as in the case of $k = 1$.

Case 3. $D \in \mathcal{N}^C$. If $D \in C \cdot X$ or $D \in C \cdot Y$, then $D \in C \cdot \mathcal{N}$ and is covered by Case 2. Otherwise, dividing by a constant, we can assume that $D = X + (a + ib)Y$, where $a, b \in R$. Thus

$$(4.9) \quad f_{\alpha\beta}(t) = g_{\alpha\beta}(t)/2\pi(i(\alpha\lambda^t + a\beta\lambda^{-t}) - b\beta\lambda^{-t}).$$

Since it suffices to prove L^2 -summability for the parts of $f_{\alpha\beta}$ corresponding to the real and imaginary parts of $g_{\alpha\beta}$ separately, we can assume $g_{\alpha\beta}$ is real in (4.9). Suppose also $b \neq 0$, since Case 2 would apply if b were 0. We note that $|f_{\alpha\beta}(t)| \leq |g_{\alpha\beta}(t)/2\pi b\beta\lambda^{-t}|$. The methods of Case 2 can be applied to prove $\sum f_{\alpha\beta} \in L^2$.

If $U \in \mathcal{U}(\mathcal{N})$, then $[D, U] = 0$, so that $\sum U f_{\alpha\beta} \in L^2$. We need to prove $\sum T^k f_{\alpha\beta} \in L^2$ for each fixed $k \in N$.

We begin with $k = 1$. Then $Tf_{\alpha\beta}(t)$ is the derivative of the right side of (4.9). In the numerator, we get various derivatives of $g_{\alpha\beta}$, while the modulus of the denominator exceeds $4\pi^2 b^2 \beta^2 \lambda^{-2t}$. The λ^{-2t} can be moved to the numerator as a derivation, and $\sum Tf_{\alpha\beta} \in L^2$. For $k > 1$, similar reasoning applies.

Case 4. $D = T + i(aX + bY)$, $a, b \in R$. Up to isomorphism, we could assume that $a = b = 1$, except for the case in which either $a = 0$ or $b = 0$. Since $T + iX$ and $T + iY$ are very similar, we need to treat only the cases $T + iX$ and $T + i(X + Y)$. Since $T + i(X + Y)$ is more complicated, we will treat this case in detail, providing brief remarks to cover the simpler case of $T + iX$.

Actually, to simplify the constants we suppose $D = T + i(X + Y)/2\pi$, so that

$$f_{\alpha\beta}(t) = \left(\int_0^t g_{\alpha\beta}(x) \exp(-(\alpha\lambda^x + \beta\lambda^{-x})/\ln \lambda) dx + C \right) \exp((\alpha\lambda^t + \beta\lambda^{-t})/\ln \lambda).$$

Subcase 4a. $\alpha > 0$ and $\beta > 0$. In this case $\exp((\alpha\lambda^t + \beta\lambda^{-t})/\ln \lambda) \rightarrow \infty$ as $t \rightarrow \pm\infty$ (or in case $b = 0$, $\exp(\alpha\lambda^t/\ln \lambda) \rightarrow \infty$ as $t \rightarrow \infty$ and $\rightarrow 1$ as $t \rightarrow -\infty$). Thus

$$(4.10) \quad \begin{aligned} f_{\alpha\beta}(t) &= - \int_t^\infty g_{\alpha\beta}(x) \exp((-(\alpha\lambda^x + \beta\lambda^{-x}) + (\alpha\lambda^t + \beta\lambda^{-t}))/\ln \lambda) dx \\ &= \int_{-\infty}^t \cdots dx. \end{aligned}$$

If $b = 0$, the $\beta\lambda^{-x}$ and $\beta\lambda^{-t}$ terms do not appear in (4.10) and we simply use $t_{\alpha\beta} = 0$ in what follows. Also, the restriction $\beta > 0$ is not necessary when we deal with the $b = 0$ case.

Note that $\frac{d}{dx}(\alpha\lambda^x + \beta\lambda^{-x}) = (\alpha\lambda^x - \beta\lambda^{-x}) \ln \lambda = 0$ iff $x = \frac{1}{2} \log_\lambda(\beta/\alpha)$, which we denote henceforth by $-1 \leq t_{\alpha\beta} \leq 0$. Since α and $\beta > 0$, $\alpha\lambda^x + \beta\lambda^{-x}$ is a decreasing function on $(-\infty, t_{\alpha\beta})$ and an increasing function on $(t_{\alpha\beta}, \infty)$. By using the *first* integral in (4.10) whenever $t \geq t_{\alpha\beta}$ and the *second* integral whenever $t < t_{\alpha\beta}$, we can assure that the exponential function in the integrand remains bounded by 1. In either case, $|g_{\alpha\beta}|$ can be bounded by one of its (X or Y) derivatives times an exponential function, with the result that $\sum f_{\alpha\beta} \in L^2$.

of $(-\infty, t_{\alpha\beta}) \cup [t_{\alpha\beta}, \infty) = (-\infty, \infty)$. (If $b = 0$, use $(X^2 + Y^2)g$ instead.) Restriction to $\lambda^{-2} \leq \beta/\alpha \leq 1$ will keep α^{-1} bounded in absolute value.

Next, let $U = (2\pi)^{-n-m} X^m Y^n \in \mathcal{U}(\mathcal{N})$. If $m - n > 0$ and $t > t_{\alpha\beta}$, then

$$|U f_{\alpha\beta}(t)| \leq \left| \alpha^m \beta^n \lambda^{(m-n)t} \int_t^\infty |g_{\alpha\beta}(x)| dx \right| \leq \int_t^\infty |U g_{\alpha\beta}(x)| dx,$$

which provides the necessary estimate. If $m - n < 0$ and $t > t_{\alpha\beta}$,

$$|U f_{\alpha\beta}(t)| \leq |\alpha^m \beta^n| \lambda^{(m-n)t} \int_t^\infty |g_{\alpha\beta}(x)| dx.$$

But $\lambda^{-2} \leq \beta/\alpha \leq 1$ implies $t_{\alpha\beta} \geq -1$, so

$$\begin{aligned} |U f_{\alpha\beta}(t)| &\leq \lambda^{n-m} \int_t^\infty |\alpha^m \beta^n g_{\alpha\beta}(x)| dx \leq \lambda^{n-m} \int_t^\infty |\alpha^{m+n} g_{\alpha\beta}(x)| dx \\ &= \lambda^{n-m} \int_t^\infty |X^{m+n} g_{\alpha\beta}(x)| \lambda^{-(m+n)x} (2\pi)^{-m-n} dx \\ &\leq \lambda^{2n} (2\pi)^{-m-n} \int_t^\infty |X^{m+n} g_{\alpha\beta}(x)| dx. \end{aligned}$$

From here the L^2 -estimates proceed as earlier.

Next, suppose $m - n < 0$ and $t < t_{\alpha\beta}$. Then $\lambda^{(m-n)t} > \lambda^{(m-n)t_{\alpha\beta}}$. Thus

$$\begin{aligned} |U f_{\alpha\beta}(t)| &\leq |\alpha^m \beta^n| \lambda^{(m-n)t} \int_{-\infty}^t |g_{\alpha\beta}(x)| dx \leq \int_{-\infty}^t |\alpha^m \beta^n \lambda^{(m-n)x} g_{\alpha\beta}(x)| dx \\ &= \int_{-\infty}^t |U g_{\alpha\beta}(x)| dx. \end{aligned}$$

The rest is as before.

Finally, for $m - n \geq 0$ and $t < t_{\alpha\beta}$, we write

$$\begin{aligned} |U f_{\alpha\beta}(t)| &\leq |\alpha^m \beta^n| \lambda^{(m-n)t} \int_{-\infty}^t |g_{\alpha\beta}(x)| dx \\ &\leq |\lambda^{2m} \beta^m \beta^n| \lambda^{(m-n)t_{\alpha\beta}} \int_{-\infty}^t |g_{\alpha\beta}(x)| dx \\ &\leq \lambda^{2m} \int_{-\infty}^t |\beta^{m+n} g_{\alpha\beta}(x)| dx \\ &= \int_{-\infty}^t |Y^{m+n} g_{\alpha\beta}(x)| \lambda^{(m+n)x} / (2\pi)^{m+n} dx \\ &\leq (2\pi)^{-m-n} \int_{-\infty}^t |Y^{m+n} g_{\alpha\beta}(x)| dx. \end{aligned}$$

The L^2 -estimates can be completed as before.

Next, we show that $\sum_{(\alpha, \beta)} T^k f_{\alpha\beta} \in L^2$, for each $k \in N$. This follows from the next lemma.

(4.11) **Lemma.** *Let f be a solution of the equation*

$$(T + i(X + Y))f = g.$$

Then $T^k f$, $k = 1, 2, 3, \dots$, is a linear combination of monomials $X^j Y^l f$ with $j + l \leq k$ plus a linear combination of X, Y, T -derivatives of g .

Proof. We proceed by induction. For $k = 1$ we have $Tf = g - iXf - iYf$. Next, $T^{k+1}f = T(T^k f) = T(X^j Y^l f)$ with $j + l \leq k$, where wlog we may assume $T^k f$ is a monomial $X^j Y^l f$.

$$\begin{aligned} T(X^j Y^l f) &= [T, X^j Y^l]f + Y^j Y^l T f \\ &= \left(\sum_{p=1}^j X \cdots X [T, X] X \cdots X Y^l + \sum_{q=1}^l X^j Y \cdots Y [T, Y] Y \cdots Y \right) f \\ &\quad + Y^j Y^l g - iX^{j+1} Y^l f - iX^j Y^{l+1} f \\ &= (j \ln \lambda X^{j+1} Y^l - l \ln \lambda X^j Y^{l+1}) f + \dots \end{aligned}$$

which is the desired expression for $T^k f$.

Remark. Similarly, $T^k f$ is a linear combination of monomials $X^j f$ with $j \leq k$ plus a linear combination of X, T -derivatives of g if f is a solution of $(T + iX)f = g$.

Subcase 4b. $\alpha < 0$ and $\beta > 0$. (The case $\alpha > 0$ and $\beta < 0$ can be treated similarly.) Once again, we have

$$f_{\alpha\beta}(t) = \left(\int_0^t g_{\alpha\beta}(x) \exp(-(\alpha\lambda^x + \beta\lambda^{-x})/\ln \lambda) dx + C \right) \exp((\alpha\lambda^t + \beta\lambda^{-t})/\ln \lambda)$$

where the terms $\beta\lambda^{-x}$ and $\beta\lambda^{-t}$ are not present if $b = 0$. Moreover, the restriction $\beta > 0$ is not needed if $b = 0$. We observe $\exp((\alpha\lambda^t + \beta\lambda^{-t})/\ln \lambda) \rightarrow 0$ as $t \rightarrow +\infty$ and $\rightarrow +\infty$ as $t \rightarrow -\infty$ (or $\rightarrow 1$ as $t \rightarrow -\infty$ in case $b = 0$). In either case, since $\lim_{t \rightarrow -\infty} f_{\alpha\beta}(t) = 0$,

$$C = \int_{-\infty}^0 g_{\alpha\beta}(x) \exp(-(\alpha\lambda^x + \beta\lambda^{-x})/\ln \lambda) dx$$

and

$$(4.12) \quad f_{\alpha\beta}(t) = \int_{-\infty}^t g_{\alpha\beta}(x) e^{\psi(x,t)} dx,$$

where

$$\psi(x, t) = (\alpha(\lambda^t - \lambda^x) + \beta(\lambda^{-t} - \lambda^{-x}))/\ln \lambda,$$

again with no $\beta(\lambda^{-t} - \lambda^{-x})$ term in case $b = 0$. We notice that $\psi(x, t) < 0$ for $x < t$. We have the estimates

$$\begin{aligned} |f_{\alpha\beta}(t)| &\leq \int_{-\infty}^t |g_{\alpha\beta}(x)| dx \leq \int_{-\infty}^t |Y g_{\alpha\beta}(x)| / |2\pi\beta\lambda^{-x}| dx \\ (4.13) \quad &\leq C \sum_{k=0}^1 \|T^k Y g_{\alpha\beta}\|_2 |\beta|^{-1} \int_{-\infty}^t \lambda^x dx \quad (\text{by Sobolev}) \\ &\leq M \sum_{k=0}^1 \|T^k Y g_{\alpha\beta}\|_2 \lambda^t / \ln \lambda \in L^2(-\infty, 0) \end{aligned}$$

since $|\beta|^{-1}$ is bounded. In fact, (4.13) implies that $\sum_{(\alpha, \beta)} f_{\alpha\beta} \in L^2(-\infty, 0)$. Next, we must consider convergence in $L^2(0, \infty)$. Thus for $t > 0$ we write

$$\begin{aligned} f_{\alpha\beta}(t) &= \int_{-\infty}^0 \cdots + \int_0^{t/2} \cdots + \int_{t/2}^t \cdots \\ &= \text{I}_{\alpha\beta}(t) + \text{II}_{\alpha\beta}(t) + \text{III}_{\alpha\beta}(t) \end{aligned}$$

where the integrands are as in (4.12).

To estimate $\text{I}_{\alpha\beta}$ we notice that for $x \leq 0 \leq t$ we have $\beta(\lambda^{-t} - \lambda^{-x}) \leq 0$ and $\alpha(\lambda^t - \lambda^x) < \alpha(\lambda^t - 1) < 0$. Hence

$$(4.14) \quad |\text{I}_{\alpha\beta}(t)| \leq e^{\alpha(\lambda^t - 1)/\ln \lambda} \int_{-\infty}^0 |g_{\alpha\beta}(x)| dx$$

and $\sum_{(\alpha, \beta)} \text{I}_{\alpha\beta} \in L^2(0, \infty)$. This is because $\alpha \gg 0$ makes the functions

$$t \mapsto \exp(\alpha(\lambda^t - \lambda^x)/\ln \lambda), \quad (\alpha, \beta) \in (\Gamma \backslash S_h)^\wedge,$$

uniformly $L^2(0, \infty)$, while the integral $\int_{-\infty}^0 \cdots$ in (4.14) can be estimated as in (4.13) making the sum finite.

For $\text{II}_{\alpha\beta}$ we have the estimate

$$\begin{aligned} (4.15) \quad |\text{II}_{\alpha\beta}(t)| &\leq \int_0^{t/2} |g_{\alpha\beta}(x)| e^{\alpha(\lambda^t - \lambda^x)/\ln \lambda} dx \\ &\leq \|g_{\alpha\beta}\|_\infty \frac{t}{2} e^{\alpha(\lambda^t - \lambda^{t/2})/\ln \lambda}. \end{aligned}$$

The right-hand side again is α -uniformly in $L^2(0, \infty)$ with $\|g_{\alpha\beta}\|_\infty$ being (α, β) -summable.

Finally,

$$\begin{aligned} (4.16) \quad |\text{III}_{\alpha\beta}(t)| &\leq \int_{t/2}^t |g_{\alpha\beta}(x)| dx \\ &\leq \int_{t/2}^t |X^m g_{\alpha\beta}(x)| / |2\pi\alpha\lambda^x|^m dx \\ &\leq \|X^m g_{\alpha\beta}\|_\infty \frac{t}{2} M^m \lambda^{-mt}, \end{aligned}$$

where M is an upper bound on $|\alpha|^{-1}$, $t\lambda^{-mt} \in L^2(0, \infty)$, and $\|X^m g_{\alpha\beta}\|_\infty$ is (α, β) -summable.

Next, we must show $\sum_{(\alpha, \beta)} U f_{\alpha\beta} \in L^2(R)$ for every fixed $U \in \mathcal{U}(\mathcal{S}_h)$.

If $U = Y^k$ we have the estimate

$$(4.17) \quad |Y^k f_{\alpha\beta}(t)| \leq \int_{-\infty}^t |Y^k g_{\alpha\beta}(x)| e^{\psi(x, t)} dx.$$

As in the beginning of Subcase 4b we can show that $\sum Y^k f_{\alpha\beta} \in L^2(R)$.

For $U = X^k$, $\sum X^k f_{\alpha\beta} \in L^2(-\infty, 0)$ because for $t \leq 0$,

$$(4.18) \quad |X^k f_{\alpha\beta}(t)| = \lambda^{2kt} |\alpha/\beta|^k |Y^k f_{\alpha\beta}(t)| \leq \lambda^2 |Y^k f_{\alpha\beta}(t)|$$

if we choose (α, β) such that $\lambda^{-2} \leq |\beta/\alpha| \leq 1$. If $t > 0$, we consider $X^k \text{I}_{\alpha\beta}$, $X^k \text{II}_{\alpha\beta}$, and $X^k \text{III}_{\alpha\beta}$ and we get the estimates (4.14), (4.15), and (4.16), each multiplied by λ^{2kt} and with $g_{\alpha\beta}$ replaced by $Y^k g_{\alpha\beta}$, as it was done in (4.18).

Finally, let $U = X^p Y^q T^r$. Case of $r \geq 1$ reduces to $r = 0$ by the Lemma (4.11). If $r = 0$, we apply X^p to $\int_{-\infty}^t |Y^q g_{\alpha\beta}(x)| e^{\psi(x,t)} dx$ as we applied $U = X^k$ to $\int_{-\infty}^t g_{\alpha\beta}(x) e^{\psi(x,t)} dx$.

Subcase 4c. $\alpha < 0$ and $\beta < 0$. We have

$$(4.19) \quad f_{\alpha\beta}(t) = \left(\int_0^t g_{\alpha\beta}(x) e^{-(\alpha\lambda^x + \beta\lambda^{-x})/\ln\lambda} dx + C \right) e^{(\alpha\lambda^t + \beta\lambda^{-t})/\ln\lambda}.$$

The function $C e^{(\alpha\lambda^t + \beta\lambda^{-t})/\ln\lambda}$ is in $\mathcal{SS}(R)$ if $\alpha < 0$ and $\beta < 0$. Hence if there is a constant C such that $f_{\alpha\beta}$ in (4.19) is super-Schwartz, then $f_{\alpha\beta} \in \mathcal{SS}(R)$ for any fixed C . We will pick $C = C_{\alpha\beta} = \int_{t_{\alpha\beta}}^0 g_{\alpha\beta}(x) e^{-(\alpha\lambda^x + \beta\lambda^{-x})/\ln\lambda} dx$, where $t_{\alpha\beta} := \frac{1}{2} \log_\lambda(\beta/\alpha)$, and α, β are chosen so that $\lambda^{-2} < \beta/\alpha \leq 1$. Thus we will work with

$$(4.19') \quad f_{\alpha\beta}(t) = \left(\int_{t_{\alpha\beta}}^t g_{\alpha\beta}(x) e^{-(\alpha\lambda^x + \beta\lambda^{-x})/\ln\lambda} dx \right) e^{(\alpha\lambda^t + \beta\lambda^{-t})/\ln\lambda}.$$

We have the estimate

$$(4.20) \quad \begin{aligned} |f_{\alpha\beta}(t)| &\leq e^{(\alpha\lambda^t + \beta\lambda^{-t})/\ln\lambda} \int_{t_{\alpha\beta}}^t (2\pi)^{-m} |(X+Y)^m g_{\alpha\beta}(x)| \\ &\quad \times e^{-(\alpha\lambda^x + \beta\lambda^{-x})/\ln\lambda} (-\alpha\lambda^x - \beta\lambda^{-x})^{-m} dx \\ &\leq e^{\dots/\ln\lambda} \|(X+Y)^m g_{\alpha\beta}\|_\infty e^{-\dots/\ln\lambda} (-2\pi \dots)^{-m} |t - t_{\alpha\beta}| \\ &\leq c \sum_{l=0}^1 \|T^l (X+Y)^m g_{\alpha\beta}\|_2 (-\alpha\lambda^t - \beta\lambda^{-t})^{-m} |t - t_{\alpha\beta}| \end{aligned}$$

where \dots stands for $\alpha\lambda^t + \beta\lambda^{-t}$. We write the inequality (4.20) for (α, β) such that

$$(4.21) \quad 2(\alpha\beta)^{1/2}/\ln\lambda > m,$$

because then the function $u \mapsto e^{-(\alpha u + \beta u^{-1})/\ln\lambda} (-\alpha u - \beta u^{-1})^{-m}$ with $u = \lambda^t$ is increasing for $t \geq t_{\alpha\beta}$ and decreasing for $t \leq t_{\alpha\beta}$. (4.21) is valid for all but a finite number of $(\alpha, \beta) \in (\Gamma \backslash S_h)^\infty$. Similarly, for $Y^l X^k f_{\alpha\beta}$ we have the estimate

$$(4.22) \quad |Y^l X^k f_{\alpha\beta}(t)| \leq c_1 \sum_{p=0}^1 \|T^p (X+Y)^m g_{\alpha\beta}\|_2 |\beta\lambda^{-t}|^l |\alpha\lambda^t|^k (-\alpha\lambda^t - \beta\lambda^{-t})^{-m} |t - t_{\alpha\beta}|.$$

But for $t \geq t_{\alpha\beta}$, if $m > k + l$

$$\begin{aligned} |\beta\lambda^{-t}|^l |\alpha\lambda^t|^k (-\alpha\lambda^t - \beta\lambda^{-t})^{-m} |t - t_{\alpha\beta}| &\leq (\beta/\alpha)^l (-\alpha)^{k+l-m} \lambda^{t(k-l-m)} |t - t_{\alpha\beta}| \\ &\leq M^{m-k-l} |t - t_{\alpha\beta}| \lambda^{t(k-l-m)} \in L^2(0, \infty) \end{aligned}$$

since $\beta/\alpha \leq 1$ and $\alpha \gg 0$.

Similarly, for $t \leq t_{\alpha\beta}$

$$\begin{aligned} |\beta\lambda^{-t}|^l |\alpha\lambda^t|^k (-\alpha\lambda^t - \beta\lambda^{-t})^{-m} |t - t_{\alpha\beta}| &\leq (\alpha/\beta)^k (-\beta)^{k+l-m} \lambda^{t(m-l+k)} |t_{\alpha\beta} - t| \\ &\leq \lambda^{2k} M^{m-k-l} |t - t_{\alpha\beta}| \lambda^{t(m-l+k)} \in L^2(-\infty, 0) \end{aligned}$$

since $\lambda^{-2} \leq \beta/\alpha$ and $\beta \gg 0$. Thus $\sum Y^l X^k f_{\alpha\beta} \in L^2(R)$. Finally,

$$\sum Y^l X^k T^m f_{\alpha\beta} \in L^2(R)$$

by Lemma (4.11).

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